

## 3D PIC SIMULATIONS OF RELATIVISTIC JETS WITH A TOROIDAL MAGNETIC FIELD

Ioana Dutan<sup>1</sup>, Ken-Ichi Nishikawa<sup>2</sup>, Athina Meli<sup>3</sup>, Oleh Kobzar<sup>4</sup>, Christoph Kohn<sup>5</sup>, Yosuke Mizuno<sup>6</sup>, Nicholas MacDonald<sup>7</sup>, Jose L. Gomez<sup>8</sup>, and Kouichi Hirota<sup>9</sup>

<sup>1</sup>Institute of Space Science, Magurele, Romania, <sup>2</sup>Alabama A&M University, Huntsville, United States, <sup>3</sup>North Carolina A&T State University, Greensboro, United States, <sup>4</sup>Institute of Nuclear Physics, Polish Academy of Sciences, Krakow, Poland, <sup>5</sup>Technical University of Denmark, National Space Institute, Kgs Lyngby, Denmark, <sup>6</sup>Tsing-Dao Lee Institute, Shanghai Jiao Tong University, Shanghai, China, <sup>7</sup>Max Planck Institute for Radio Astronomy, Bonn, Germany, <sup>8</sup>Instituto de Astrofísica de Andalucía CSIC, Granada, Spain, <sup>9</sup>Taiwan Institute of Astronomy and Astrophysics, Academia Sinica, Taipei, China

**Abstract.** The properties of relativistic jets, their interaction with the ambient environment, and particle acceleration due to kinetic instabilities are studied self-consistently with Particle-in-Cell (PIC) simulations. An important key issue is how a toroidal magnetic field affects the evolution of an electron-positron and an electron-ion jet, how kinetic instabilities such as the Weibel instability (WI), the mushroom instability (MI) and the kinetic Kelvin-Helmholtz instability (kKH) are excited, and how such instabilities contribute to particle acceleration comparing simulations without a toroidal magnetic field (Nishikawa et al. 2016). We show that WI, MI and kKH excited at the linear stage, generate a quasi-steady  $x$ -component of an electric field which accelerates and decelerates electrons. In this work, we use a new jet injection scheme where an electric current is self-consistently generated at the jet orifice by the jet particles. We inject both electron-positron and electron-proton jets with a toroidal magnetic field (with a top-hat jet density profile) for a sufficiently long time in order to examine the non-linear effects of the jet evolution. We observe significant differences in the structure of the strong electromagnetic fields that are driven by the kinetic instabilities. We find that different jet compositions present different strongly excited instability modes. The magnetic field in the non-linear stage generated by different instabilities becomes dissipated and reorganized into a new topology. The 3-dimensional magnetic field topology indicates possible reconnection sites and the accelerated particles are significantly accelerated in the non-linear stage by the dissipation of the magnetic field and/or reconnection. This study will shed further light on the nature of astrophysical relativistic magnetized jet phenomena.

### 1. Introduction

Relativistic jets are ubiquitous in astrophysical systems, such as active galactic nuclei and gamma ray bursts.

Understanding jet composition, particle acceleration, and emission of radiation is still an ongoing work among many research groups (including Nishikawa et al. since 2009) with PIC simulations of global jets containing helical magnetic fields (Nishikawa et al. 2016, 2020)

#### 1.1. Key scientific questions

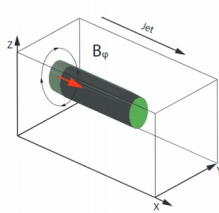
- How do global jets evolve with different particle species?
- How do toroidal magnetic fields affect the kinetic instabilities in plasma jets and their nonlinear evolution, as well as the magnetic reconnection?
- Jets in Jets really happen due to magnetic reconnection

#### 1.2. Why do we need to perform PIC simulations of relativistic jets

- Kinetic instabilities (e.g., kinetic Kelvin-Helmholtz Instability (kKH), Mushroom Instability (MI), and Weibel Instability (WI)) are a key issue in understanding jet evolution besides the kink instability in Relativistic Magnetohydrodynamic (RMHD) simulations
- Helical (toroidal) magnetic fields are crucial in understanding these instabilities
- PIC global jet simulations are a new and innovative approach which provides insights of the complex evolution of relativistic jets with kinetic processes including emission of radiation which cannot be performed by RMHD simulations
- Nonthermal acceleration due to reconnection may generate flares

### 2. Simulation set-up

- Develop a new jet injection scheme: a cylindrical jet is injected into the ambient plasma at rest, which propagates in the  $x$ -direction, with a toroidal magnetic field. Plasma jet is composed of electrons and positrons and of electrons and ions (with  $m_i/m_e = 4$ ), respectively
- The jet is injected at  $x = 100\Delta$  in the center of the  $y$ - $z$  plane at  $(y_c, z_c)$ , and propagates in the  $x$ -direction
- Jet radial width in cylindrical coordinates is  $r_{jet} = 100\Delta$
- Equations of toroidal magnetic field ( $B_\phi$ ) followed from Mizuno et al. (2014) and Nishikawa et al. (2020), where the components of  $B_\phi$  are written in Cartesian coordinates ( $B_x$  and  $B_y$ ), and  $B_z = 0$
- Equation for the current  $J$ , follows from  $\partial E/\partial t = \text{curl } B = J$

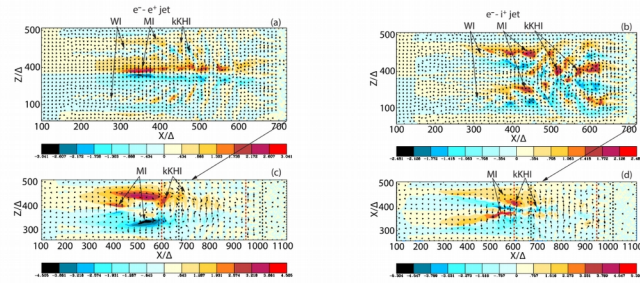


The schematic jet injection scheme with a toroidal magnetic field ( $B_\phi$ ). The jet electrons and positrons (ions) are injected so that the current (indicated with the red arrow) is generated to support the toroidal magnetic field.

- To sustain the current in the jet, the toroidal magnetic field is gradually applied at the jet orifice ( $\omega\Delta = 1.00 - 1.02$ ), and a magnetic electric field is set up  $E_{\text{ext}} = -\nabla \phi$ ,  $\Delta$  is the grid cell size and  $v_x = v_{jx}$ , where  $v_{jx}$  is the  $x$ -component of the jet velocity

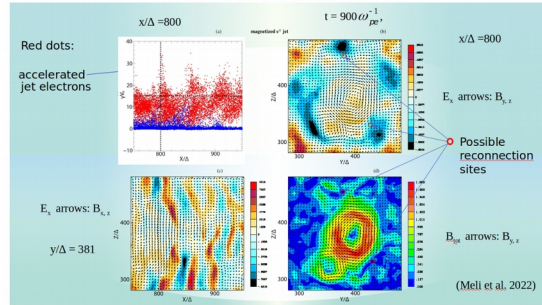
### 3. Results (Meli et al. 2022, in preparation)

#### 3.1. Generation of kinetic instabilities and global evolution of the jets containing toroidal magnetic fields

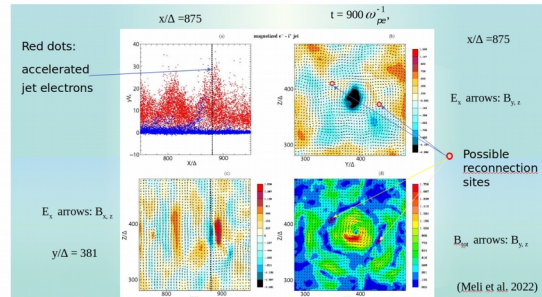


Color maps of the magnetic field amplitude  $B_y$  and arrows depicting the magnetic field components in the  $x$ - $z$  plane, both at  $t = 600\omega_{pe}^{-1}$  (upper panels) and  $900\omega_{pe}^{-1}$  (lower panels), respectively. The jet is injected at  $x = 100\Delta$  in the middle of the  $y$ - $z$  plane and propagates in  $+x$ -direction. Panels (a, c) are for an  $e^+e^-$  plasma while panels (b, d) are for an  $e^-e^-$  composition.

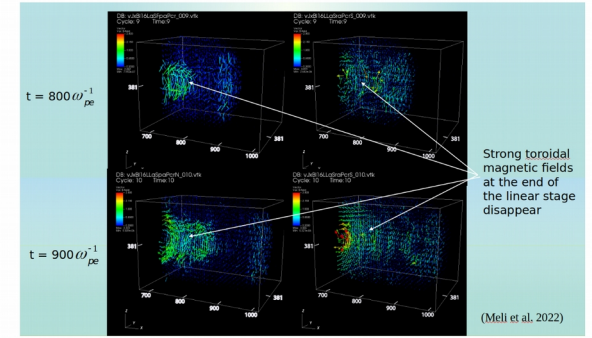
#### 3.2. Electromagnetic fields and particle acceleration



For an  $e^+$  plasma jet: Panel a:  $x$ - $yy$ -phase-space distribution of jet (red) and ambient (blue) electrons. Panels b, c, & d: structures generated due to different MI modes. Panels b & c: possible magnetic reconnection sites.



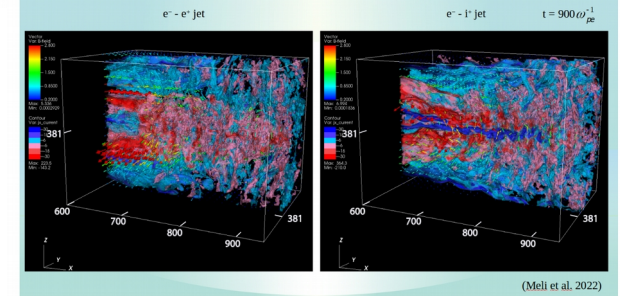
For an  $e^-e^-$  plasma jet: Panel a: The  $x$ - $yy$ -phase-space distribution of jet (red) and ambient (blue) electrons. Panels b, c, & d: structures generated due to different MI modes. Panels b & c: possible magnetic reconnection sites.



Strong toroidal magnetic fields at the end of the linear stage disappear

(Meli et al. 2022)

Evolution of the magnetic field near the jet head: Left-hand panels for an  $e^+$  plasma field and right-hand panels for an  $e^-e^-$  composition, acceleration of the electrons in the non linear stage, between  $800 < x/\Delta < 1000$  is not due to the electric field of the generated instabilities, but is correlated to the dissipation of the magnetic fields.



The  $J$  structure and magnetic fields. Color map: Structure of the current along jet direction of propagation,  $J$ . Thin lines: Magnetic field. Left panel shows the growth of the kKH and MI (and WI), as well as the generation of two modes of MI with twisted magnetic fields in the region  $750 < x/\Delta < 950$  which shows the non-linear saturation at  $x/\Delta < 950$  indicating the non-linear saturation of the grown instabilities. Right panel shows the strong negative current in the center of the jet and the toroidal magnetic field which is opposite to the original direction.

### 4. Summary

- We simulated electron-positron and electron-ion relativistic jets containing toroidal magnetic fields
- For an electron-positron jet a MI is excited and combined with a kKH while a quasi steady  $E_x$  modulates jet particles
- For an electron-ion jet, kinetic instabilities are dominated by the MI
- MI and kKH produce quasi-steady electric field ( $E_x$ ) for both types of plasma jets
- These electric fields accelerate and decelerate electrons and ions
- Electrons are further accelerated due to turbulent magnetic fields generated by dissipations of helical magnetic fields (reconnection)
- Further investigations is important in order to confirm and/or find other acceleration mechanisms with varying simulation parameters such as, jet radius, magnetization factor, jet structure (Gaussian shape), etc.

**Bibliography:** Nishikawa, K.-I., Nemeic, J., Meebey, M., et al., 2009, *ApJ*, 698, L10  
Nishikawa, K.-I., Mizuno, Y., Nemeic, J., et al., 2016, *Galaxies*, 4, 38  
Nishikawa, K.-I., Mizuno, Y., Gomez, J. L., et al., 2020, *MNRAS*, 491, 2022

#### Acknowledgments

The authors would like to thank the collaborators Martin Pohl and Jack Nemeic for the insightful discussions during the development of this work. This work was supported by the NASA-NSX2-MERG, NNX14AP-21G, and NNX14AP-64G grants. Recent work was also provided by the NASA through the Cluster Award NNX17-011J (PI: Meng-Gan Li) issued by the Cluster & Sun Center, which is operated by the SSO for and on behalf of the NASA under contract NAS8-01060. The work of J.N. has been supported by Nanoscale Common Nuclei through research project 2019/13 B57W0260. Y.M. is supported by the ERC Synergy Grant "BlackHoles: Can Images the Event Horizon of Black Holes" (Grant No. 810139). The work of I.D. has been supported by the NSFC 51875402. Simulations were performed using resources and software facilities at NASA Advanced Supercomputing (NAS) using Comet at the San Diego Supercomputing Center, and Hubs at the Pittsburgh Supercomputing Center, which are supported by the NSF. I.D. acknowledges the support of the Spanish Ministerio de Economía y Competitividad (grants AYA2019-108999-BI00-00-C21), the Conselleria de Economia, Comercio, Emprego y I+D+i of the Xunta de Galicia (grant PI18-110), the Conselleria Superior de Investigación Científica (grant 2019/AS/12), and the State Agency for Research of the Spanish MCIU through the Center of Excellence Severo Ochoa award for the Instituto de Astrofísica de Andalucía (SEV-2017-0709).